

A global Carleman estimate in a transmission wave equation and application to a one-measurement inverse problem

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Abstract

We consider a transmission wave equation in two embedded domains in \mathbb{R}^2 , where the speed is $a_1 > 0$ in the inner domain and $a_2 > 0$ in the outer domain. We prove a global Carleman inequality for this problem under the hypothesis that the inner domain is strictly convex and $a_1 > a_2$. As a consequence of this inequality, uniqueness and Lipschitz stability are obtained for the inverse problem of retrieving a stationary potential for the wave equation with Dirichlet data and discontinuous principal coefficient from a single time-dependent Neumann boundary measurement.

1. Introduction and main results

1.1. Presentation of the problem

The inverse problem of recovering coefficients from a wave equation with discontinuous coefficients from boundary measurements arises naturally in geophysics and, more precisely, in seismic prospection of Earth inner layers [20].

Here, we are interested in the case where only one particular measurement is available. This could be important, for instance, in seismic prospection, where data of a single wave that propagates through the Earth are considered.

Consider two embedded domains, where the speed coefficients are $a_1 > 0$ in the inner domain and $a_2 > 0$ in the outer domain. Stability of the inverse problem we study here is

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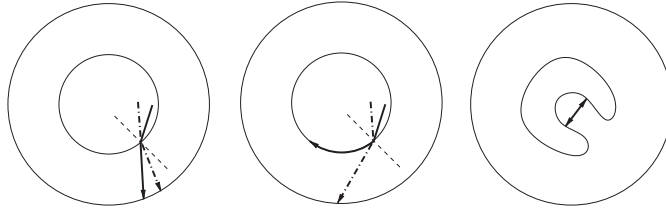


Figure 1. Two domains with speed coefficients a_1 (inner) and a_2 (outer). In the first figure (left), if $a_1 > a_2$ by Snell's law all the inner rays reach the exterior boundary independently of their incident angles. Conversely, in the second figure (centre) if $a_1 < a_2$ some rays with large incident angles remain trapped near the inner interface. The last figure (right) shows a trapped ray into a captive domain.

obtained by deriving a global Carleman estimate for the wave equation with discontinuous coefficients. We prove this Carleman inequality in the case the inner domain is strictly convex and the speed is monotonically increasing from the outer to the inner layers, i.e. $a_1 > a_2$. This last situation is, incidentally, the general case into the Earth.

Figure 1 illustrates the role of these hypotheses and gives some intuition with the help of Snell's law. In the case $a_1 > a_2$ (see figure 1, left) the incident rays coming from the inner domain towards the outer domain become closer to the normal at the interface since $\sin(\theta_1) > \sin(\theta_2)$, where $\theta_i, i = 1, 2$, are the corresponding incident angles. Therefore, all the rays coming from the inner ball with any incident angle θ_1 in $(-\pi/2, \pi/2)$ succeed in crossing the interface. In the opposite case, when $a_1 < a_2$ (see figure 1, centre) we have $\sin(\theta_1) < \sin(\theta_2)$ and there is a critical incident angle $\theta^* < \pi/2$ such that the rays with incident angles θ_1 out of the range $(-\theta^*, \theta^*)$ remain supported near the interface and do not reach the outer domain, so this information does not arrive at the exterior boundary. Finally, strict convexity of the inner domain avoids trapped rays (see figure 1, right).

Global Carleman estimates and the method of Bukhgeim–Klibanov [7, 9] are especially useful for solving the one-measurement inverse problems. It is possible to obtain local Lipschitz stability around the single known solution, provided that this solution is regular enough and contains enough information [27] (see also [26] and [41]). Many other related inverse results for hyperbolic equations use the same strategy. A complete list is too long to be given here. To cite some of them see [36] and [41] where Dirichlet boundary data and Neumann measurements are considered and [23, 24] where Neumann boundary data and Dirichlet measurements are studied. These references are all based upon the use of local or global Carleman estimates. Related to this, there are also general pointwise Carleman estimates that are also useful in similar inverse problems [16, 17, 28].

Recently, global Carleman estimates and applications to one-measurement inverse problems were obtained in the case of variable but still regular coefficients, see [22] for the isotropic case, and [30] and [4] for the anisotropic case. It is interesting to note that these authors require a bound on the gradient of the coefficients, so that the idea of approximating discontinuous coefficients by smooth ones is not useful.

There are a number of important works [8, 19, 37–39] concerning the same inverse problem in the case that several boundary measurements are available. In these cases, it is possible to retrieve speed coefficients and even discontinuity interfaces without any restrictive hypothesis of strict convexity or speed monotonicity. For instance, one can retrieve the interface by observing the traveltimes reflection of several waves. Indeed, it is well known that the interface can be recovered as the envelope of certain curves as shown in figure 2 (see also [20] and the references therein). This method works independently of the sign of $a_1 - a_2$ and

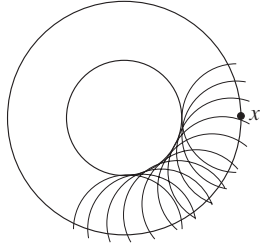


Figure 2. Recovering the interface as the envelope of circumferences. Each circumference, centred at some point x on the exterior boundary, represents the possible locations of the nearest point of the inner interface where the reflection took place for a given traveltime measured at x .

this explains in part why there are no geometrical or speed monotonic hypotheses for these kinds of inverse results.

Let us now give some insight into the relationship between this work and exact controllability or energy decay for the wave equation with discontinuous coefficients.

Firstly, the global Carleman estimate we obtain immediately implies a particular case of a well-known result of exact controllability for the transmission wave equation [31]. Roughly speaking, the result of [31] states that we can control internal waves from the exterior boundary in a layered speed media if the speed is monotonically increasing from the outer to the inner layers and the inner domain is star shaped, a weaker assumption than strict convexity. Moreover, if the speed monotonicity is inverted, there are non-controllable solutions with concentrated energy near the interface [12, 33].

Secondly, there exist several results about the growth of the resolvent for the spectral stationary transmission problem, from where it is possible to derive the speed of local energy decay for the evolution wave equation with transmission conditions [10]. In the case $a_1 > a_2$ and if the inner domain is strictly convex, it has been shown using micro-local analysis [34, 35] that the speed of the energy decay is exponential if the dimension of space is odd and polynomial otherwise. In the general case, including the cases when $a_1 < a_2$ or the inner domain is not strictly convex, it has been proved using micro-local analysis and global Carleman estimates for the spectral problem [3] that the energy decays as the inverse of the logarithm of the time.

Note that we shall only consider here the case of a discontinuous coefficient which is constant on each subdomain (i.e. a_1 and a_2 constants). We will indeed concentrate our discussion on the main difficulty, namely the discontinuity at the interface. However, we could also consider variable coefficients $a_1(x)$ and $a_2(x)$ such that their traces at the interface are constant, under additional assumptions of boundedness of ∇a_j similar to those appearing in [22] (see remark 2.1.5).

Finally, we note that a global Carleman estimate [15] has also been obtained for the heat equation with discontinuous coefficients. That work was initially motivated by the study of the exact null controllability of the semilinear wave equation, but the estimate has been recently used to prove local Lipschitz stability for a one-measurement inverse problem for the heat equation with discontinuous coefficients [5, 6].

Having introduced the problem, let us now present our main results.

1.2. Inverse problem

Let Ω and Ω_1 be two open subsets of \mathbb{R}^2 with smooth boundaries Γ and Γ_1 . Suppose that Ω_1 is simply connected, $\overline{\Omega_1} \subset \Omega$ and set $\Omega_2 = \Omega \setminus \overline{\Omega_1}$. Thus, we have $\partial\Omega_2 = \Gamma \cup \Gamma_1$. We also set

$$a(x) = \begin{cases} a_1 & x \in \Omega_1 \\ a_2 & x \in \Omega_2 \end{cases} \quad (1)$$

with $a_j > 0$ for $j = 1, 2$. We consider the following wave equation:

$$\begin{cases} u_{tt} - \operatorname{div}(a(x)\nabla u) + p(x)u = 0 & (x, t) \in \Omega \times (0, T) \\ u = 0 & (x, t) \in \Gamma \times (0, T) \\ u(0) = u_0 & x \in \Omega \\ u_t(0) = u_1 & x \in \Omega. \end{cases} \quad (2)$$

We know that [13, 32] for each $p \in L^\infty(\Omega)$, $u_0 \in H_0^1(\Omega)$ and $u_1 \in L^2(\Omega)$, there exists a unique weak solution $u(p)$ of equation (2) such that

$$u(p) \in C([0, T]; H_0^1(\Omega)), \quad u_t \in C([0, T]; L^2(\Omega)).$$

We shall prove the well-posedness of the inverse problem consisting of retrieving the potential p involved in equation (2), by knowing the flux (the normal derivative) of the solution $u(p)$ of (2) on the boundary. We will prove uniqueness and stability of the nonlinear inverse problem characterized by the nonlinear application:

$$p|_\Omega \mapsto a_2 \frac{\partial u}{\partial \nu} \Big|_{\Gamma \times (0, T)}. \quad (3)$$

More precisely, we will answer the following questions:

Uniqueness. Does the equality $\frac{\partial u(q)}{\partial \nu} = \frac{\partial u(p)}{\partial \nu}$ on $\Gamma \times (0, T)$ imply $q = p$ on Ω ?

Stability. Is it possible to estimate $(q - p)|_\Omega$ by $\left(\frac{\partial u(q)}{\partial \nu} - \frac{\partial u(p)}{\partial \nu}\right)|_{\Gamma \times (0, T)}$ in suitable norms?

The idea is to reduce the *nonlinear inverse problem* to some *perturbed inverse problem* which will be solved with the help of a global Carleman estimate. More precisely, we will give a local answer about the determination of p , working first on the perturbed version of the problem, as shown in section 3. Assuming that $p \in L^\infty$ is a given function, we are concerned with the stability around p . That is to say, p and $u(p)$ are known while q and $u(q)$ are unknown.

We are able to prove the following result, which states the stability of the inverse problem.

Theorem 1. *Assume Ω is bounded, $\Omega_1 \subset\subset \Omega$ is a strictly convex domain with boundary Γ_1 of class C^3 and $a_1 > a_2 > 0$. There exists $T_0 > 0$ such that, given $T > T_0$, if $p \in L^\infty(\Omega)$, $u_0 \in H_0^1(\Omega)$, $u_1 \in L^2(\Omega)$ and $r > 0$ satisfy*

- $|u_0(x)| \geq r > 0$ a.e. in Ω and
- $u(p) \in H^1(0, T; L^\infty(\Omega))$,

then, given a bounded set $\mathcal{U} \subset L^\infty(\Omega)$, there exists a constant

$C = C(a_1, a_2, \Omega_1, \Omega_2, T, \|p\|_{L^\infty(\Omega)}, \|u(p)\|_{H^1(0, T; L^\infty)}, \mathcal{U}, \nabla) > \iota$ such that

$$\|p - q\|_{L^2(\Omega)} \leq C \left\| a_2 \frac{\partial u(p)}{\partial \nu} - a_2 \frac{\partial u(q)}{\partial \nu} \right\|_{H^1(0, T; L^2(\Gamma))}$$

for all $q \in \mathcal{U}$, where $u(p)$ and $u(q)$ are the solutions of (2) with potential p and q , respectively.

Remark 1. In section 3, an estimate for T_0 as functions of a_1, a_2, Ω_1 and Ω_2 is given. See theorem 3.

Let us remark that as a direct consequence of the *local stability* of theorem 1 we have the following *global uniqueness* for our inverse problem:

Corollary 1. *If $u(p)$ and $u(q)$ are two solutions of (2) for potentials p and q in $L^\infty(\Omega)$ with $u(p), u_0, u_1, a_1, a_2, \Omega_1, \Omega_2$ and T satisfying the hypothesis of theorem 1 and such that $\frac{\partial u(q)}{\partial \nu} = \frac{\partial u(p)}{\partial \nu}$ on $\Gamma \times (0, T)$ then $p = q$.*

The proof of this result consists of two parts: a global Carleman estimate and the resolution of the inverse problem and Lipschitz stability following the methods introduced in [9] and [27] which we have already mentioned in the introduction.

1.3. Carleman estimate

We introduce here our main result concerning a global Carleman estimate for the solutions of problem (2) extended to the time interval $(-T, T)$. We set $Q = \Omega \times (-T, T)$, $\Sigma = \Gamma \times (-T, T)$, $\Sigma_1 = \Gamma_1 \times (-T, T)$, $Q_j = \Omega_j \times (-T, T)$, $u_j = u \mathbb{1}_{Q_j}$ and ν_j the outward unit normal to Ω_j , for $j = 1, 2$.

We will work with an equivalent formulation of (2). Note that for each $f \in L^2(Q)$, u solves the equation

$$u_{tt} - \operatorname{div}(a \nabla u) + pu = f \quad \text{in } Q \quad (4)$$

if and only if, for each $j \in \{1, 2\}$, u_j solves (see [31])

$$u_{j,tt} - a_j \Delta u_j + pu_j = f \mathbb{1}_{Q_j} \quad \text{in } Q_j \quad (5)$$

together with the *transmission conditions*

$$\begin{cases} u_1 = u_2 & \text{on } \Sigma_1 \\ a_1 \frac{\partial u_1}{\partial \nu_1} + a_2 \frac{\partial u_2}{\partial \nu_2} = 0 & \text{on } \Sigma_1. \end{cases} \quad (6)$$

In order to construct a convenient weight function, take $x_0 \in \Omega_1$ and for each $x \in \Omega \setminus \{x_0\}$ define $\ell(x_0, x) = \{x_0 + \lambda(x - x_0) : \lambda \geq 0\}$. Since Ω_1 is convex there is exactly one point $y(x)$ such that

$$y(x) \in \Gamma_1 \cap \ell(x_0, x). \quad (7)$$

We define the function $\rho : \Omega \setminus \{x_0\} \rightarrow \mathbb{R}^+$ by

$$\rho(x) = |x_0 - y(x)|. \quad (8)$$

Let $\varepsilon > 0$ be such that $\overline{B_\varepsilon} \subset \Omega_1$ (and small enough in a sense that we will precise later) and let $0 < \varepsilon_1 < \varepsilon_2 < \varepsilon$. Then, we consider a cut-off function $\eta \in C^\infty(\mathbb{R})$ such that

$$0 \leq \eta \leq 1, \quad \eta = 0 \quad \text{in } B_{\varepsilon_1}(x_0), \quad \eta = 1 \quad \text{in } \Omega \setminus \overline{B_{\varepsilon_2}}(x_0). \quad (9)$$

For each $j \in \{1, 2\}$ we take k such that $\{j, k\} = \{1, 2\}$ and we define the following functions in the whole domain $\Omega \times \mathbb{R}$:

$$\phi_j(x, t) = \eta(x) \frac{a_k}{\rho(x)^2} |x - x_0|^2 - \beta t^2 + M_j \quad (x, t) \in \Omega \times \mathbb{R}, \quad (10)$$

where β, M_1 and M_2 are positive numbers that will be chosen later. Then, the weight function we will use in this work is

$$\phi(x, t) = \begin{cases} \phi_1(x, t) & (x, t) \in \Omega_1 \times \mathbb{R} \\ \phi_2(x, t) & (x, t) \in \Omega_2 \times \mathbb{R}. \end{cases} \quad (11)$$

Note that (see (c) and (d) in proposition 1) ϕ_1 and ϕ_2 satisfy (6) if and only if

$$M_1 - M_2 = a_1 - a_2. \quad (12)$$

We denote

$$L = \partial_t^2 - a\Delta \quad \text{and} \quad E(z) = |z_t|^2 - a|\nabla z|^2.$$

As usual, we do the change of variables

$$\varphi = e^{\lambda\phi}, \lambda > 0, \quad w = e^{s\varphi}u, s > 0, \quad P(w) = e^{s\varphi}L(e^{-s\varphi}w) \quad (13)$$

and after algebraical computations we split $P(w)$ into three terms as follows:

$$P(w) = P_1(w) + P_2(w) + R(w),$$

where for some fixed real number $\gamma \in (0, 1)$

$$\begin{aligned} P_1(w) &= w_{tt} - a\Delta w + s^2\lambda^2\varphi^2 E(\phi)w, \\ P_2(w) &= (\gamma - 1)s\lambda\varphi L(\phi)w - s\lambda^2\varphi E(\phi)w - 2s\lambda\varphi(\phi_t w_t - a\nabla\phi \cdot \nabla w), \\ R(w) &= -\gamma s\lambda\varphi L(\phi)w. \end{aligned}$$

We will write $P^\phi, P_1^\phi, P_2^\phi$, etc if we want to make the dependence on ϕ explicit. Also, given $U \subset \mathbb{R}^2$, we define the norm in $H^1(U \times (-T, T))$, given by

$$\|g\|_{U,\varphi}^2 = s\lambda \int_{-T}^T \int_U (|g_t|^2 + |\nabla g|^2)\varphi + s^3\lambda^3 \int_{-T}^T \int_U |g|^2\varphi^3, \quad (14)$$

and

$$\Sigma_+^\phi = \{(x, t) \in \Sigma : \nabla\phi(x, t) \cdot \nu(x) > 0\}. \quad (15)$$

Finally, we define the space

$$\begin{aligned} X &= \{u \in L^2(-T, T; L^2(\Omega)) \mid Lu_j \in L^2(-T, T; L^2(\Omega_j)), j = 1, 2; u|_\Sigma = 0, \\ &\quad u(\pm T) = u_t(\pm T) = 0, \text{ and } u \text{ satisfies (6)}\}. \end{aligned}$$

The main global Carleman estimate is the following:

Theorem 2. *Assume Ω_1 is a strictly convex domain of class C^3 , and $a_1 > a_2 > 0$. Let $x_k \in \Omega_1, k = 1, 2$ and let ϕ^k, φ^k, w^k be the corresponding functions defined for x_k as we did before for x_0 in (7), (8), (11) and (13). Let ν be the unit outward normal to Ω . Then there exist $C > 0, s_0 > 0$ and $\lambda_0 > 0$ such that*

$$\begin{aligned} &\sum_{k=1}^2 (\|P_1^{\phi^k}(w^k)\|_{L^2(Q)}^2 + \|P_2^{\phi^k}(w^k)\|_{L^2(Q)}^2 + \|w^k\|_{Q,\varphi^k}^2) \\ &\leq C \sum_{k=1}^2 \left(\|P^{\phi^k}(w^k)\|_{L^2(Q)}^2 + s\lambda \int \int_{\Sigma_+^{\phi^k}} \varphi^k \left| a_2 \frac{\partial w^k}{\partial \nu} \right|^2 \right) \end{aligned} \quad (16)$$

for all $u \in X, \lambda \geq \lambda_0$ and $s \geq s_0$.

Note that on the right-hand side of (16) we have the term

$$\|P^{\phi^k}(w^k)\|_{L^2(Q)}^2 = \int \int_Q e^{2s\varphi^k} |Lu|^2.$$

Since we consider an equation given by the operator $L_p = \partial_{tt} - a\Delta + p$, it is important to note that the same estimate with the right-hand side equal to $\int \int_Q e^{2s\varphi^k} |L_p u|^2$ is also true for all potentials p such that $|p|_{L^\infty(\Omega)} \leq m$, with m already fixed. Indeed,

$$|Lu|^2 \leq 2|L_p u|^2 + 2m^2|u|^2$$

and taking s large enough, the left-hand side of the Carleman estimate of theorem 2 can absorb the term $2Cm^2 \int \int_Q e^{2s\varphi} |u|^2$. That is, we have the following result.

Corollary 2. *Under the hypothesis and notation of theorem 2, given $m \in \mathbb{R}$, there exist $C > 0$ (depending on m), $s_0 > 0$ and $\lambda_0 > 0$ such that for all $p \in L^\infty(\Omega)$ with $\|p\|_{L^\infty(\Omega)} \leq m$ we have*

$$\begin{aligned} & \sum_{k=1}^2 (\|P_1^{\phi^k}(w^k)\|_{L^2(Q)}^2 + \|P_2^{\phi^k}(w^k)\|_{L^2(Q)}^2 + \|w^k\|_{Q, \phi^k}^2) \\ & \leq C \sum_{k=1}^2 \left(\iint_Q e^{2s\phi^k} |L_p u|^2 + s\lambda \iint_{\Sigma_+^{\phi^k}} \varphi^k \left| a_2 \frac{\partial w^k}{\partial \nu} \right|^2 \right) \end{aligned} \quad (17)$$

for all $u \in X$, $\lambda \geq \lambda_0$ and $s \geq s_0$.

The paper is organized as follows. In section 2, we give the proof of theorem 2. In section 3, we apply inequality (17) to derive the stability of the inverse problem presented in theorem 1. The proof of theorem 2 is organized in several subsections. In proposition 1, we prove the properties of ϕ which will allow us to use it as a weight in a Carleman estimate. In subsection 2.2, with ϕ as the weight function, we develop the L^2 -product $\langle P_1(w), P_2(w) \rangle$ for functions with non-zero boundary values. We prove inequality (16) in subsection (2.3).

2. Proof of the Carleman inequality

2.1. Weight function

Here, we prove that the function ϕ satisfies enough properties for being a weight function in a Carleman estimate.

We will use the following notation:

$$\begin{aligned} M &= M_1 \mathbb{1}_{Q_1} + M_2 \mathbb{1}_{Q_2}, & \bar{a} &= a_2 \mathbb{1}_{Q_1} + a_1 \mathbb{1}_{Q_2}, & c(x) &= \frac{\bar{a}}{\rho(x)^2}, \\ \Omega_0 &= \Omega_1 \cup \Omega_2, & \Omega_{x_0} &= \Omega_0 \setminus \bar{B}_\varepsilon(x_0), \\ Q_0 &= \Omega_0 \times (-T, T), & Q_{x_0} &= \Omega_{x_0} \times (-T, T). \end{aligned}$$

Proposition 1. *If Ω_1 is a strictly convex domain of class C^3 , we can take $\varepsilon, \delta > 0$ such that*

- (a) $|\nabla \phi| \geq \delta > 0$ in $Q_{x_0} = (\Omega_1 \cup \Omega_2) \setminus B_\varepsilon(x_0) \times (-T, T)$,
- (b) $\nabla \phi_1(x, t) \cdot \nu_1(x) \geq \delta > 0$, $\forall (x, t) \in \Sigma_1$,

where ν_1 is the unit outward normal vector to Ω_1 .

If additionally (12) is satisfied, we also have

- (c) $\phi_1(x, t) = \phi_2(x, t) = a_2 - \beta t^2 + M_1 \forall (x, t) \in \Sigma_1$,
- (d) $a_1 \frac{\partial^{\alpha_1 + \alpha_2}}{\partial x_1^{\alpha_1} \partial x_2^{\alpha_2}} \phi_1(x, t) = a_2 \frac{\partial^{\alpha_1 + \alpha_2}}{\partial x_1^{\alpha_1} \partial x_2^{\alpha_2}} \phi_2(x, t)$ for all $(x, t) \in \Sigma_1$ and $\alpha_1, \alpha_2 \in \mathbb{N} \cup \{0\}$ with $\alpha_1 + \alpha_2 \leq 3$,
- (e) $\Delta \phi(x, t) \geq 2c(x) \quad \forall (x, t) \in Q_{x_0}$,
- (f) $D^2(\phi)(X, X) \geq \delta_1 |X|^2$ in $Q_{x_0} \forall X \in \mathbb{R}^2$ for some $\delta_1 > 0$.

Proof. We have

$$\nabla \phi = 2c(x)(x - x_0) + |x - x_0|^2 \nabla c(x).$$

By definition, $\rho(x)$ (thus also $c(x)$) is constant in the direction of $x - x_0$. Therefore

$$(x - x_0) \cdot \nabla c(x) = 0$$

and

$$\begin{aligned} |\nabla\phi|^2 &= 4c^2(x)|x - x_0|^2 + |x - x_0|^4|\nabla c(x)|^2 \\ &\geq 4c^2(x)|x - x_0|^2 \\ &\geq 4\left(\frac{\bar{a}}{\text{diam}(\Omega)^2}\right)^2 \varepsilon^2 \quad \text{in } \Omega_{x_0} \end{aligned}$$

and (a) is proved.

Now, it is clear that $\phi_1(x, t) = a_2 - \beta t^2 + M_1$ for each $(x, t) \in \Sigma_1$, so $\Gamma_1 \times \{t\}$ is a level curve of $\phi_1(\cdot, t)$ for each $t \in [-T, T]$. Since $\phi_1(x, t) < a_2 - \beta t^2 + M_1 < \phi_1(y, t)$ for any $x \in \Omega_1$ and $y \in \Omega_2$, we have $\nabla\phi_1 = |\nabla\phi_1|v_1$ on Σ_1 and thus (a) implies (b).

By definition $\rho(x) = |x - x_0|$ for all $x \in \Gamma_1$, hence (c) is simply deduced from (12).

Without loss of generality, we can take $x_0 = 0$. Writing ρ in polar coordinates, Γ_1 can be parameterized by

$$\gamma(\theta) = (\rho(\theta) \cos \theta, \rho(\theta) \sin \theta), \quad (18)$$

and then ρ is a C^3 function. If $D = \frac{\partial^{\alpha_1 + \alpha_2}}{\partial x_1^{\alpha_1} \partial x_2^{\alpha_2}}$ with $\alpha_1 + \alpha_2 \leq 3$, we get

$$a_1 D\phi_1(x, t) = a_1 a_2 D\left(\frac{|x - x_0|^2}{\rho(x)^2}\right) = a_2 D\phi_2(x, t)$$

for all $(x, t) \in \Sigma_1$ and (d) is proved.

The expression for the Hessian matrix of second derivatives in polar coordinates is

$$D^2(\phi) = Q_\theta H(\phi) Q_\theta^T$$

where Q_θ is the rotation matrix by angle θ and

$$H(\phi) = \begin{pmatrix} \frac{\partial^2 \phi}{\partial r^2} & \frac{1}{r} \left(\frac{\partial^2 \phi}{\partial r \partial \theta} - \frac{1}{r} \frac{\partial \phi}{\partial \theta} \right) \\ \frac{1}{r} \left(\frac{\partial^2 \phi}{\partial r \partial \theta} - \frac{1}{r} \frac{\partial \phi}{\partial \theta} \right) & \frac{1}{r^2} \frac{\partial^2 \phi}{\partial \theta^2} + \frac{1}{r} \frac{\partial \phi}{\partial r} \end{pmatrix}. \quad (19)$$

Now, we have that (recall that $x_0 = 0$)

$$\phi(\theta, r, t) = \frac{\bar{a}}{\rho(\theta)^2} r^2 - \beta t^2 + M.$$

One can note that ϕ is well defined and smooth in Ω_{x_0} (which means $\{r \geq \varepsilon\} \setminus \Gamma_1$). All the computations that follow are valid in this set. We have already said above that ρ is constant with respect to r and only depends on θ such that $\frac{\partial \rho}{\partial r} = 0$. Hence, we have that

$$H(\phi) = \frac{2\bar{a}}{\rho^2} \begin{pmatrix} 1 & -\frac{\rho_\theta}{\rho} \\ -\frac{\rho_\theta}{\rho} & \frac{1}{\rho^2} (3\rho_\theta^2 - \rho\rho_{\theta\theta} + \rho^2) \end{pmatrix}, \quad (20)$$

where we have denoted $\rho_\theta = \frac{\partial \rho}{\partial \theta}$ and so on.

We will use the following well-known facts (see, for example, [18]) concerning curves in the plane:

Lemma 1. *Let γ be a C^2 curve in the plane. Then,*

(a) *the curve γ is strictly convex only at those points where their curvature is positive,*

(b) if γ is parameterized in polar coordinates by its angle, that is

$$\gamma(\theta) = (r(\theta) \cos \theta, r(\theta) \sin \theta),$$

then the curvature of γ is given by the formula

$$\kappa_\gamma(\theta) = \frac{r^2 + 2r_\theta^2 - rr_{\theta\theta}}{(r^2 + r_\theta^2)^{3/2}}.$$

Since the polar parameterization of Γ_1 is precisely the above function with $r(\theta) = \rho(\theta)$ and Ω_1 is strictly convex, we obtain

$$\kappa_{\Gamma_1}(\theta) = \frac{\rho^2 + 2\rho_\theta^2 - \rho\rho_{\theta\theta}}{(\rho^2 + \rho_\theta^2)^{3/2}} > 0 \quad \forall \theta \in [0, 2\pi[. \quad (21)$$

From (20) and (21), we have

$$\begin{aligned} \Delta\phi &= \text{tr}(D^2(\phi)) \\ &= \text{tr}(H(\phi)) \\ &= 2c \left(1 + \frac{1}{\rho^2} (3\rho_\theta^2 - \rho\rho_{\theta\theta} + \rho^2) \right) \\ &\geq 2c \left(1 + \frac{1}{\rho^2} (\rho^2 + \rho_\theta^2)^{3/2} \kappa_{\Gamma_1} \right) > 2c \end{aligned}$$

and (e) is proved.

We also have

$$\begin{aligned} \det(H(\phi)) &= \frac{2c}{\rho^2} (3\rho_\theta^2 - \rho\rho_{\theta\theta} + \rho^2 - \rho_\theta^2) \\ &= \frac{2c}{\rho^2} (\rho^2 + \rho_\theta^2)^{3/2} \kappa_{\Gamma_1} > 0. \end{aligned}$$

By the Sylvester's Criterion we can see that $H(\phi)$ (thus $D^2(\phi)$) is positive definite. Indeed, the element $\frac{2c}{\rho^2}$ and the determinant of the matrix $H(\phi)$ are positive. Finally, since $\bar{\Omega}$ is compact, this implies (f) and the proof of lemma 1 is complete. \square

We introduce now the last hypothesis we will need in order to get the Carleman inequality:

$$\beta < \min \left\{ \frac{\min\{a_1, a_2\}\delta_1}{2}, \frac{M}{T^2} \right\} \quad (22)$$

and we take $\gamma \in (0, 1)$ such that

$$\gamma > \frac{2\beta}{\beta + \frac{a_1 a_2}{\text{diam}(\Omega)^2}} \quad (23)$$

$$\gamma < \frac{2 \min\{a_1, a_2\}\delta_1}{2\beta + \max\{a_1, a_2\} \|\Delta\phi\|_{L^\infty(\Omega_0)}^2}. \quad (24)$$

Remark 2.

- (i) We have to take M large enough in order to (22) and the hypothesis of the inverse problem (see theorem 3) become compatible.
- (ii) Taking β small enough, (23) and (24) become compatible. But, as we will see in the next section, the smaller is the β the bigger need to be the inversion time for the inverse problem (see theorem 3).

- (iii) Actually, the optimal δ_1 is the first eigenvalue of $D^2(\phi)$. Having an explicit expression for it could help to a better choice of β .
- (iv) From the hypothesis of theorem 2 we have $a_2 < a_1$ and the maximum and the minimum in (22) and (24) are known.

Remark 3. *Pseudoconvexity.* It is easy to check that proposition 1 and assumption (22) imply that ϕ_1 and ϕ_2 are pseudoconvex [1, 21] with respect to P in $\Omega_1 \setminus B_\varepsilon(x_0)$ and Ω_2 , respectively. Global Carleman estimates (without explicit dependence on the parameter λ) can be deduced in each subdomain $\Omega_1 \setminus B_\varepsilon(x_0)$ and Ω_2 , for solutions that vanish on the exterior boundary and the interface (see [1, 40]). Nevertheless, the traces of the transmission wave equation do not vanish on the interface Γ_1 and we will need to carefully use hypothesis $a_1 > a_2$ and the parameter s . On the other hand, we have to use the parameter λ in order to get rid of the lack of pseudoconvexity in $B_\varepsilon(x_0)$.

Remark 4. Coefficient $a = a(x)$ variable. We can also prove a Carleman estimate in the more general case $a(x) = a_1(x)\mathbb{1}_{\Omega_1} + a_2(x)\mathbb{1}_{\Omega_2}$ with $a_j \in C^1(\overline{\Omega}_j)$, $j = 1, 2$, if each a_j is constant in the interface Γ_1 , and under some hypotheses on ∇a (similar to those of [22]). More precisely, if we check the pseudoconvexity condition in this case, we will have that ϕ is pseudoconvex with respect to the operator $\partial_{tt} - \operatorname{div}(a\nabla u)$ (in each domain Q_j) if there exists $\theta \in (0, 1)$ such that

$$\frac{|\nabla\phi \cdot \nabla a_j|}{2a_j} < \delta_1(1 - \theta), \quad j = 1, 2 \quad (25)$$

$$\beta < \frac{\sqrt{a_j}}{\sqrt{a_j} + T|\nabla a_j|} \left(\frac{a_j \delta_1 \theta}{2} \right), \quad j = 1, 2. \quad (26)$$

It is easy to check that these hypotheses are compatible with the assumptions of theorem 1. Indeed, for T sufficiently large, there exists β satisfying both $\beta > a_1/T^2$ and (26). Nevertheless, in order to construct the weight function as we have done above (and deal with the traces of the solutions on the interface Γ_1), it is crucial for each function a_j to be constant on the interface.

2.2. Listing all the terms

In this part of the work we develop the L^2 -product of $P_1(w)$ and $P_2(w)$. We will do formal computations, by writing generically ϕ for the weight function and Q for the domain with boundary Σ .

As presented in section 1, we have, for $\lambda > 0$, $s > 0$,

$$\begin{aligned} \varphi &= e^{\lambda\phi}, & w &= e^{s\varphi}u \\ P(w) &= e^{s\varphi}L(e^{-s\varphi}w) = P_1(w) + P_2(w) + R(w) \end{aligned}$$

where

$$\begin{aligned} P_1(w) &= w_{tt} - a\Delta w + s^2\lambda^2\varphi^2E(\phi)w, \\ P_2(w) &= (\gamma - 1)s\lambda\varphi L(\phi)w - s\lambda^2\varphi E(\phi)w, & -2s\lambda\varphi(\phi_t w_t - a\nabla\phi \cdot \nabla w). \\ R(w) &= -\gamma s\lambda\varphi L(\phi)w. \end{aligned}$$

We set $\langle P_1(w), P_2(w) \rangle_{L^2} = \sum_{i,j=1}^3 I_{i,j}$, where $I_{i,j}$ is the integral of the product of the i th term in $P_1(w)$ and the j th term in $P_2(w)$. Therefore,

$$\begin{aligned}
I_{1,1} &= -s\lambda(\gamma - 1) \iint_Q \varphi L(\phi) |w_t|^2 + \frac{s\lambda^2(\gamma - 1)}{2} \iint_Q |w|^2 \varphi (\phi_{tt} + \lambda |\phi_t|^2) L(\phi), \\
I_{1,2} &= s\lambda^2 \iint_Q |w_t|^2 \varphi E(\phi) - s\lambda^2 \iint_Q |w|^2 \varphi |\phi_{tt}|^2 - \frac{5s\lambda^3}{2} \iint_Q |w|^2 \varphi \phi_{tt} |\phi_t|^2 \\
&\quad + \frac{s\lambda^3}{2} \iint_Q |w|^2 a \varphi \phi_{tt} |\nabla \phi|^2 - \frac{s\lambda^4}{2} \iint_Q |w|^2 \varphi |\phi_t|^2 E(\phi), \\
I_{1,3} &= s\lambda \iint_Q |w_t|^2 \varphi (\phi_{tt} + \lambda |\phi_t|^2) - 2s\lambda^2 \iint_Q w_t \phi_t \varphi a \nabla w \cdot \nabla \phi \\
&\quad + s\lambda \iint_Q |w_t|^2 \varphi a (\Delta \phi + \lambda |\nabla \phi|^2) - s\lambda \iint_\Sigma |w_t|^2 a \varphi \nabla \phi \cdot \nu, \\
I_{2,1} &= -s\lambda(\gamma - 1) \iint_\Sigma \varphi L(\phi) w a \frac{\partial w}{\partial \nu} + s\lambda(\gamma - 1) \iint_Q a |\nabla w|^2 \varphi L(\phi) \\
&\quad - s\lambda^2 \frac{\gamma - 1}{2} \iint_Q |w|^2 \varphi a L(\phi) (\lambda |\nabla \phi|^2 + \Delta \phi) - s\lambda \frac{\gamma - 1}{2} \iint_Q |w|^2 a \varphi \Delta(L(\phi)) \\
&\quad - s\lambda^2(\gamma - 1) \iint_Q |w|^2 a \varphi (\nabla \phi \cdot \nabla L(\phi)) + s\lambda \frac{\gamma - 1}{2} \iint_\Sigma |w|^2 a \varphi \nabla L(\phi) \cdot \nu \\
&\quad + s\lambda^2 \frac{\gamma - 1}{2} \iint_\Sigma |w|^2 a \varphi L(\phi) \frac{\partial \phi}{\partial \nu}, \\
I_{2,2} &= s\lambda^2 \iint_\Sigma a \varphi E(\phi) w \nabla w \cdot \nu - \frac{s\lambda^2}{2} \iint_\Sigma a |w|^2 (\lambda \varphi E(\phi) \nabla \phi + \varphi \nabla E(\phi)) \cdot \nu \\
&\quad + \frac{s\lambda^3}{2} \iint_Q |w|^2 a \varphi E(\phi) (\Delta \phi + \lambda |\nabla \phi|^2) - 2s\lambda^3 \iint_Q |w|^2 a^2 \varphi D^2(\phi) (\nabla \phi, \nabla \phi) \\
&\quad + \frac{s\lambda^2}{2} \iint_Q |w|^2 a \varphi \Delta(E(\phi)) - s\lambda^2 \iint_Q |\nabla w|^2 a \varphi E(\phi), \\
I_{2,3} &= s\lambda \iint_Q |\nabla w|^2 a \varphi L(\phi) + s\lambda^2 \iint_Q |\nabla w|^2 a \varphi E(\phi) + 2s\lambda^2 \iint_Q a^2 \varphi |\nabla \phi \cdot \nabla w|^2 \\
&\quad - 2s\lambda^2 \iint_Q a \varphi \phi_t w_t \nabla w \cdot \nabla \phi + 2s\lambda \iint_Q a^2 \varphi D^2(\phi) (\nabla w, \nabla w) \\
&\quad + s\lambda \iint_\Sigma |\nabla w|^2 a^2 \varphi \nabla \phi \cdot \nu + 2s\lambda \iint_\Sigma a \varphi (\phi_t w_t - a \nabla \phi \cdot \nabla w) \frac{\partial w}{\partial \nu}, \\
I_{3,1} &= s^3 \lambda^3 (\gamma - 1) \iint_Q |w|^2 \varphi^3 L(\phi) E(\phi), \\
I_{3,2} &= -s^3 \lambda^4 \iint_Q |w|^2 \varphi^3 E(\phi)^2, \\
I_{3,3} &= s^3 \lambda^3 \iint_Q |w|^2 \varphi^3 E(\phi) L(\phi) + 2s^3 \lambda^3 \iint_Q |w|^2 \varphi^3 (|\phi_t|^2 \phi_{tt} + a^2 D^2(\phi) (\nabla \phi, \nabla \phi)) \\
&\quad + 3s^3 \lambda^4 \iint_Q |w|^2 \varphi^3 E(\phi)^2 + s^3 \lambda^3 \iint_\Sigma a |w|^2 \varphi^3 E(\phi) \frac{\partial \phi}{\partial \nu}.
\end{aligned}$$

Gathering all these terms, we get

$$\begin{aligned}
\langle P_1(w), P_2(w) \rangle_{L^2(Q)} &= 2s\lambda \iint_Q |w_t|^2 \varphi \phi_{tt} - \gamma s\lambda \iint_Q |w_t|^2 \varphi L(\phi) \\
&\quad + 2s\lambda^2 \iint_Q \varphi (|w_t|^2 |\phi_t|^2 - 2w_t \phi_t a \nabla w \cdot \nabla \phi + a^2 |\nabla \phi \cdot \nabla w|^2)
\end{aligned}$$

$$\begin{aligned}
& + 2s\lambda \iint_Q a^2 \varphi D^2(\phi)(\nabla w, \nabla w) \\
& + \gamma s\lambda \iint_Q a |\nabla w|^2 \varphi L(\phi) + 2s^3 \lambda^4 \iint_Q |w|^2 \varphi^3 E(\phi)^2 \\
& + 2s^3 \lambda^3 \iint_Q |w|^2 \varphi^3 (|\phi_t|^2 \phi_{tt} + a^2 D^2(\phi)(\nabla \phi, \nabla \phi)) \\
& + \gamma s^3 \lambda^3 \iint_Q |w|^2 \varphi^3 L(\phi) E(\phi) + X + J,
\end{aligned}$$

where J is the sum of all the boundary terms:

$$\begin{aligned}
J = s\lambda \iint_\Sigma & \left(a^2 \varphi |\nabla w|^2 \frac{\partial \phi}{\partial \nu} - 2a^2 \varphi (\nabla \phi \cdot \nabla w) \frac{\partial w}{\partial \nu} \right) \\
& + s\lambda \frac{\gamma - 1}{2} \iint_\Sigma |w|^2 \varphi a \nabla L(\phi) \cdot \nu + s\lambda^2 \frac{\gamma - 1}{2} \iint_\Sigma |w|^2 \varphi a L(\phi) \frac{\partial \phi}{\partial \nu} \\
& - s\lambda(\gamma - 1) \iint_\Sigma w a \frac{\partial w}{\partial \nu} \varphi L(\phi) + s\lambda^2 \iint_\Sigma a \varphi E(\phi) w \frac{\partial w}{\partial \nu} \\
& - s\lambda^3 \frac{1}{2} \iint_\Sigma |w|^2 \varphi E(\phi) a \frac{\partial \phi}{\partial \nu} - s\lambda^2 \frac{1}{2} \iint_\Sigma |w|^2 a \varphi \nabla E(\phi) \cdot \nu \\
& + 2s\lambda \iint_\Sigma a \varphi \phi_t w_t \frac{\partial w}{\partial \nu} - s\lambda \iint_\Sigma |w_t|^2 \varphi a \frac{\partial \phi}{\partial \nu} + s^3 \lambda^3 \iint_\Sigma |w|^2 \varphi^3 E(\phi) a \frac{\partial \phi}{\partial \nu}
\end{aligned}$$

and X is the sum of the remaining terms, in such a way that

$$|X| \leq C s \lambda^3 \iint_Q \varphi^3 |w|^2$$

In the following, we denote by A_j , $j = 1, \dots, 8$, the first eight integrals we have listed in the product of $P_1(w)$ by $P_2(w)$. Thus, we have

$$\langle P_1(w), P_2(w) \rangle_{L^2(Q)} = \sum_{j=1}^8 A_j + X + J. \quad (27)$$

2.3. Proof of theorem 2

We take β , γ and M satisfying the hypotheses (12), (22), (23) and (24), and ϕ as the corresponding weight function. We assume throughout all this part of the work the hypothesis of proposition 1 (especially that Ω_1 is strictly convex with C^3 boundary).

Recall the notation $Q_0 = (\Omega_1 \cup \Omega_2) \times (-T, T)$ and $Q_{x_0} = (\Omega_0 \setminus B_\varepsilon(x_0)) \times (-T, T)$. We apply the above computations to $w = e^{s\varphi} u$ with $u \in X$ in each one of the open sets Q_1 and Q_2 . Adding the terms that result in both cases, we have (recall that $u(\pm T) = u_t(\pm T) = 0$ for all $u \in X$)

$$\langle P_1(w), P_2(w) \rangle_{L^2(Q_0)} = \sum_{j=1}^8 A_{j, Q_0}(w) + X_{Q_0}(w) + J_{\Sigma_1}(w) + J_\Sigma(w), \quad (28)$$

where we have written A_{j, Q_0} instead of the integral A_j given in subsection 2.2 taken in the set Q_0 , etc.

The proof of theorem 2 is based on the following three facts:

- The sum of the A_j -integrals in Q_{x_0} can be minored.
- The sum of terms in the interface given by J_{Σ_1} is nonnegative.

- We can introduce a second weight function centred at a point different to x_0 in order to deal with the integrals in $B_\varepsilon(x_0)$.

The key points in each step of the proof are based on the properties of ϕ listed in proposition 1.

2.3.1. The interior

Proposition 2. *There exist $\delta > 0$, $C > 0$ and $\lambda_0 > 0$ such that*

$$\sum_{j=1}^8 A_{j, Q_0}(w) \geq \delta \|w\|_{\Omega_0, \varphi} - C \|w\|_{B_\varepsilon(x_0), \varphi}$$

for all $\lambda \geq \lambda_0$ and for all $u \in X$.

Proof. We arrange the terms into four groups:

- (i) $A_{1, Q_0} + A_{2, Q_0} = s\lambda \iint_{Q_0} |w_t|^2 \varphi (-\gamma L(\phi) - 4\beta)$. For all (x, t) in Q_{x_0} we have

$$\begin{aligned} -\gamma L(\phi) - 4\beta &= \gamma(2\beta + a\Delta\phi) - 4\beta \\ &\geq \gamma(2\beta + 2ac(x)) - 4\beta && \text{by proposition 1} \\ &\geq \gamma(2\beta + 2a\frac{\bar{a}}{\text{diam}(\Omega)^2}) - 4\beta && \text{by definition of } c(x) \\ &= \delta > 0 && \text{by (23)}. \end{aligned}$$

Therefore,

$$A_{1, Q_0} + A_{2, Q_0} \geq \delta_1 s\lambda \iint_{\Omega_{x_0}} |w_t|^2 \varphi - Cs\lambda \iint_{B_\varepsilon(x_0)} |w_t|^2 \varphi.$$

- (ii) $A_{3, Q_0} = 2s\lambda^2 \iint_{Q_0} \varphi (\phi_t w_t - a\nabla\phi \cdot \nabla w)^2 \geq 0$.
 (iii) $A_{4, Q_{x_0}} + A_{5, Q_{x_0}} = s\lambda \iint_{Q_{x_0}} \varphi (a^2 D^2(\phi)(\nabla w, \nabla w) + \gamma a L(\phi) |\nabla w|^2)$. Then, using proposition 1 and (24), we obtain

$$\begin{aligned} A_{4, Q_{x_0}} + A_{5, Q_{x_0}} &\geq s\lambda \iint_{Q_{x_0}} \varphi (2a^2 \delta_1 |\nabla w|^2 + \gamma a L(\phi) |\nabla w|^2) \\ &\geq s\lambda \iint_{Q_{x_0}} \varphi a(2a\delta_1 - \gamma(2\beta + a\Delta\phi)) |\nabla w|^2 \\ &\geq s\lambda \iint_{Q_{x_0}} \varphi a(2a\delta_1 - \gamma(2\beta + a\|\Delta\phi\|_{L^\infty})) |\nabla w|^2 \\ &\geq s\lambda \delta_2 \iint_{Q_{x_0}} \varphi |\nabla w|^2. \end{aligned}$$

Therefore,

$$A_{4, Q_0} + A_{5, Q_0} \geq \delta_2 s\lambda \iint_{\Omega_{x_0}} \varphi |\nabla w|^2 - Cs\lambda \iint_{B_\varepsilon(x_0)} \varphi |\nabla w|^2.$$

- (iv) $\sum_{j=6}^8 A_{j, Q_0} = s^3 \lambda^3 \iint_{Q_0} |w|^2 \varphi^3 F_\lambda(\phi)$ where

$$\begin{aligned} F_\lambda(\phi) &= 2\lambda E(\phi)^2 + 2|\phi_t|^2 \phi_{tt} + 2a^2 D^2(\phi)(\nabla\phi, \nabla\phi) + \gamma L(\phi) E(\phi) \\ &= 2\lambda E(\phi)^2 + \gamma L(\phi) E(\phi) - 16\beta^3 t^2 + 2a^2 D^2(\phi)(\nabla\phi, \nabla\phi) \\ &= 2\lambda E(\phi)^2 + \underbrace{(\gamma L(\phi) - 4\beta) E(\phi)}_{b(x) < 0} - 4\beta a |\nabla\phi|^2 + 2a^2 D^2(\phi)(\nabla\phi, \nabla\phi). \end{aligned}$$

From proposition 1 and (22), there exists $d_0 > 0$ such that for all $(x, t) \in Q_{x_0}$ we have

$$\begin{aligned} F_\lambda(\phi) &\geq 2\lambda E(\phi)^2 + b(x)E(\phi) - 4\beta a|\nabla\phi|^2 + 2a^2\delta_1|\nabla\phi|^2 \\ &\geq 2\lambda E(\phi)^2 + b(x)E(\phi) + a(2a\delta_1 - 4\beta)|\nabla\phi|^2 \\ &\geq 2\lambda E(\phi)^2 - \|b\|_\infty|E(\phi)| + a(2a\delta_1 - 4\beta)|\nabla\phi|^2 \\ &\geq 2\lambda E(\phi)^2 - \|b\|_\infty|E(\phi)| + d_0 \\ &= \begin{cases} f_{1,\lambda}(E(\phi)) & \text{if } E(\phi) > 0 \\ f_{2,\lambda}(E(\phi)) & \text{if } E(\phi) < 0 \end{cases} \end{aligned}$$

where

$$\begin{aligned} f_{j,\lambda} &: \mathbb{R} \longrightarrow \mathbb{R} \\ x &\longmapsto 2\lambda x^2 + (-1)^j \|b\|_\infty x + d_0. \end{aligned}$$

As $d_0 > 0$, there exists $\lambda_0 > 0$ such that for all $\lambda \geq \lambda_0$

$$\min_{\mathbb{R}}(f_{j,\lambda}) \geq \frac{d_0}{2} > 0 \quad j = 1, 2.$$

Thus, for each $\lambda \geq \lambda_0$ we have

$$\sum_{j=6}^8 A_{j,Q_0} \geq \delta_3 s^3 \lambda^3 \iint_{Q_{x_0}} |w|^2 \varphi^3 - C s^3 \lambda^3 \iint_{B_\varepsilon(x_0)} |w|^2 \varphi^3.$$

By collecting all the terms A_{j,Q_0} together, we conclude the proof of proposition 2. \square

2.3.2. The interface. Since the interface Γ_1 is a common boundary of Ω_1 and Ω_2 , the term J_{Σ_1} in (28) is the sum of the integrals coming from each domain: $J_{\Sigma_1} = J_{\Sigma_1}(w_1) + J_{\Sigma_1}(w_2)$. We have the following result:

Proposition 3. *Suppose $0 < a_2 < a_1$. Then there exists $s_0 > 0$ such that*

$$J_{\Sigma_1} = J_{\Sigma_1}(w_1) + J_{\Sigma_1}(w_2) \geq 0 \quad \forall s \geq s_0$$

for all $\varepsilon \in X$.

Proof. We enumerate the ten integrals arising in (27) associated with the common boundary Σ_1 , and we denote by J_i the sum of the i th integral in (27) which comes from Ω_1 with the respective one of Ω_2 : $J_i = J_i(w_1) + J_i(w_2)$.

In order to prove the inequality, we arrange the terms into three groups. In each case, we use proposition 1 and the fact that w satisfies the transmission conditions.

- (i) It is not difficult to see that $J_k = 0$ for each $k \in \{3, 4, 7, 8, 9\}$. Indeed, from (d) of proposition 1 we get $L_1(\phi_1) = L_2(\phi_2)$ and $a_1 \nabla E_1(\phi_1) = a_2 \nabla E_2(\phi_2)$ on Σ_1 , and the desired result follows.

Now, let us denote by g the real function defined in Σ_1 as

$$g(x, t) := E_1(\phi_1) - E_2(\phi_2) = \left(\frac{1}{a_2} - \frac{1}{a_1} \right) \left| a_1 \frac{\partial \phi}{\partial \nu_1} \right|^2.$$

Since $a_2 < a_1$, we have $g > 0$ in Σ_1 .

Thus, we can prove the following:

(ii) $J_2 + J_6 + \frac{1}{2}J_{10} \geq 0 \quad \forall s \geq s_0$.

Indeed,

$$\begin{aligned} -J_2 - J_6 &= s\lambda \frac{1-\gamma}{2} \iint_{\Sigma_1} |w|^2 \varphi a \nabla L(\phi) \cdot \nu + s\lambda^3 \frac{1}{2} \iint_{\Sigma_1} |w|^2 \varphi \left(a_1 \frac{\partial \phi_1}{\partial \nu_1} \right) g(x, t) \\ &\leq \frac{1}{2} J_{10} \end{aligned}$$

for all $s \geq s_0$, since $\varphi \geq 1$.

(iii) $J_1 + J_5 + \frac{1}{2}J_{10} \geq 0 \forall s \geq s_0$.

By construction, ϕ is constant on each level $\Gamma_1 \times \{t\}$ of the interface (proposition 1).

Thus,

$$\nabla \phi_j \cdot \nabla w_j = \frac{\partial \phi_j}{\partial \nu_j} \frac{\partial w_j}{\partial \nu_j} \quad \text{in } \Sigma_1 \quad \text{for } j = 1, 2. \quad (29)$$

Moreover, since w satisfies (6) we have

$$\left| \frac{\partial w_1}{\partial \tau_1} \right| = \left| \frac{\partial w_2}{\partial \tau_2} \right| \quad \text{in } \Sigma_1 \forall u \in X.$$

Hence,

$$\sum_{j=1}^2 \left| \frac{\partial w_j}{\partial \tau_j} \right|^2 a_j^2 \varphi_j \frac{\partial \phi_j}{\partial \nu_j} = \left| \frac{\partial w_1}{\partial \tau_1} \right|^2 \varphi_1 \left(a_1 \frac{\partial \phi_1}{\partial \nu_1} \right) (a_1 - a_2) > 0. \quad (30)$$

From (29) and (30), we get

$$\begin{aligned} J_1 &\geq -s\lambda \iint_{\Sigma_1} \left(\left| a_1 \frac{\partial w_1}{\partial \nu_1} \right|^2 \varphi_1 \frac{\partial \phi_1}{\partial \nu_1} + \left| a_2 \frac{\partial w_2}{\partial \nu_2} \right|^2 \varphi_2 \frac{\partial \phi_2}{\partial \nu_2} \right) \\ &= -s\lambda \iint_{\Sigma_1} \left| a_1 \frac{\partial w_1}{\partial \nu_1} \right|^2 \varphi_1 \left(\frac{\partial \phi_1}{\partial \nu_1} + \frac{\partial \phi_2}{\partial \nu_2} \right) \\ &= s\lambda \iint_{\Sigma_1} \left| a_1 \frac{\partial w_1}{\partial \nu_1} \right|^2 \varphi_1 \left(\frac{1}{a_2} - \frac{1}{a_1} \right) \left(a_1 \frac{\partial \phi_1}{\partial \nu_1} \right). \end{aligned}$$

On the other hand,

$$\begin{aligned} -J_5 &= -s\lambda^2 \iint_{\Sigma_1} w_1 \varphi_1 \left(a_1 \frac{\partial w_1}{\partial \nu_1} \right) g(x, t) \\ &\leq \frac{1}{2} s^2 \lambda^3 \iint_{\Sigma_1} |w_1|^2 \varphi_1 g + \frac{1}{2} \lambda \iint_{\Sigma_1} \varphi_1 g \left| a_1 \frac{\partial w_1}{\partial \nu_1} \right|^2 \\ &= \frac{1}{2} s^2 \lambda^3 \iint_{\Sigma_1} |w_1|^2 \varphi_1 g + \frac{1}{2} \lambda \iint_{\Sigma_1} \left| a_1 \frac{\partial w_1}{\partial \nu_1} \right|^2 \varphi_1 \left(\frac{1}{a_2} - \frac{1}{a_1} \right) \left| a_1 \frac{\partial \phi_1}{\partial \nu_1} \right|^2 \\ &\leq \frac{1}{2\delta} s^2 \lambda^3 \iint_{\Sigma_1} |w|^2 \varphi_1 g \left(a_1 \frac{\partial \phi_1}{\partial \nu_1} \right) \\ &\quad + \frac{C}{2} \lambda \iint_{\Sigma_1} \left| a_1 \frac{\partial w_1}{\partial \nu_1} \right|^2 \varphi_1 \left(\frac{1}{a_2} - \frac{1}{a_1} \right) \left(a_1 \frac{\partial \phi_1}{\partial \nu_1} \right) \\ &\leq \frac{1}{2} J_{10} + J_1 \quad \forall s \geq s_1. \end{aligned}$$

Proposition (3) is proved. \square

2.3.3. *The boundary Σ .* Since we deal with functions w such that $w_2 = 0$ in Σ , we have

$$\begin{aligned}
J_\Sigma &= J_1(w_2) = -s\lambda \iint_\Sigma \varphi \left| a_2 \frac{\partial w}{\partial \nu} \right|^2 \left(\frac{\partial \phi}{\partial \nu} \right) \\
&\geq -s\lambda \iint_{\Sigma_+} \varphi \left| a_2 \frac{\partial w}{\partial \nu} \right|^2 \left(\frac{\partial \phi}{\partial \nu} \right) \\
&\geq -s\lambda \iint_{\Sigma_+} \varphi \left| a_2 \frac{\partial w}{\partial \nu} \right|^2 \left\| \frac{\partial \phi}{\partial \nu} \right\|_{L^\infty(\Sigma)} \\
&= -s\lambda C \iint_{\Sigma_+} \varphi \left| a_2 \frac{\partial w}{\partial \nu} \right|^2,
\end{aligned} \tag{31}$$

where we have defined $\Sigma_+ = \{(x, t) \in \Gamma : \nabla \phi(x, t) \cdot \nu(x) > 0\}$.

2.3.4. *Carrying all together.* From (28), (31) and propositions 2 and 3, there exist s_0 , λ_0 , $C \in \mathbb{R}$ such that for each $s \geq s_0$ and $\lambda \geq \lambda_0$ we have

$$\|w\|_{\Omega_0, \varphi}^2 - C \|w\|_{B_\varepsilon(x_0), \varphi}^2 + X_{Q_0} - s\lambda C \iint_{\Sigma_+} \varphi \left| a_2 \frac{\partial w}{\partial \nu} \right|^2 \leq C \langle P_1(w), P_2(w) \rangle_{L^2(Q_0)}. \tag{32}$$

Adding $\frac{C}{2} (|P_1(w)|_{L^2(Q_0)}^2 + |P_2(w)|_{L^2(Q_0)}^2)$ to both sides of (32) we obtain

$$\begin{aligned}
&|P_1(w)|_{L^2(Q_0)}^2 + |P_2(w)|_{L^2(Q_0)}^2 + \|w\|_{\Omega_0, \varphi}^2 - C \|w\|_{B_\varepsilon(x_0), \varphi}^2 \\
&\quad + X_{Q_0} - s\lambda C \iint_{\Sigma_+} \varphi \left| a_2 \frac{\partial w}{\partial \nu} \right|^2 \leq C |P_1(w) + P_2(w)|_{L^2(Q_0)}^2 \\
&= C |P(w) - R(w)|_{L^2(Q_0)}^2 \\
&\leq C (|P(w)|_{L^2(Q_0)}^2 + |R(w)|_{L^2(Q_0)}^2).
\end{aligned} \tag{33}$$

Thanks to (22) we have $\lambda \leq C\varphi$ for λ large enough. Therefore,

$$\begin{aligned}
|X_{Q_0}| + C |R(w)|_{L^2(Q_0)}^2 &\leq Cs\lambda^3 \iint_Q \varphi^2 |w|^2 \quad \forall \lambda \geq \lambda_1 \\
&\leq \frac{1}{2} \|w\|_{\Omega_0, \varphi}^2 \quad \forall s \geq s_1.
\end{aligned} \tag{34}$$

From (33) and (34) we get, for all $s \geq \max\{s_0, s_1\}$, $\lambda \geq \max\{\lambda_0, \lambda_1\}$,

$$\begin{aligned}
&|P_1(w)|_{L^2(Q_0)}^2 + |P_2(w)|_{L^2(Q_0)}^2 + \|w\|_{\Omega_0, \varphi}^2 \leq C |P(w)|_{L^2(Q_0)}^2 \\
&\quad + C \|w\|_{B_\varepsilon(x_0), \varphi}^2 + s\lambda C \iint_{\Sigma_+} \varphi \left| a_2 \frac{\partial w}{\partial \nu} \right|^2.
\end{aligned} \tag{35}$$

2.3.5. *Eliminating the term in $B_\varepsilon(x_0)$.* In the last step we will remove the integral in $B_\varepsilon(x_0)$ from the right-hand side of (35). In order to do that, first remark that x_0 can be arbitrarily chosen in Ω_1 since Ω_1 is strictly convex.

Thus, we can take two different points in Ω_1 and we have the two respective inequalities given by (35). Now, we will show that the left-hand side of each inequality can absorb the term $\|\cdot\|_{B_\varepsilon(x_0)}$ from the other inequality provided that ε is small and λ is large enough.

Denote by x_1, x_2 two points in Ω_1 , and by ϕ^1, ϕ^2 their respective weight functions. In order to have $\|\cdot\|_{B_\varepsilon(x_1), \phi^1}$ absorbed by the term $\|w^2\|_{\Omega_0, \varphi^2}$ it suffices that

$$C\phi^1 < \frac{1}{2}\phi^2 \quad \text{in } B_\varepsilon(x_1),$$

i.e.,

$$e^{\lambda(\phi^2 - \phi^1)} > 2C \quad \text{in } B_\varepsilon(x_1).$$

Thus, if we show that it is possible to have $\phi^2 - \phi^1 > \delta > 0$ in $B_\varepsilon(x_1)$ by taking λ large enough we are done.

In fact, let $d = \frac{1}{2}|x_1 - x_2|$ and assume that $\varepsilon < d$. Then, for all $x \in B_\varepsilon(x_1)$ we have

$$\begin{aligned} \phi^1(x, t) &\leq \frac{a}{\rho_1^2} \varepsilon^2 - \beta t^2 + M \\ &\leq \frac{a}{\alpha_1^2} \varepsilon^2 - \beta t^2 + M, \end{aligned} \quad (36)$$

where $\alpha_1 = d(x_1, \Gamma_1) > 0$.

In the same way, if we denote $D_2 = \max_{y \in \Gamma_1} d(y, x_2)$, we get

$$\begin{aligned} \phi^2(x, t) &\geq \frac{a}{\rho_2^2} d^2 - \beta t^2 + M \\ &\geq \frac{a}{D_2^2} d^2 - \beta t^2 + M \quad \forall x \in B_\varepsilon(x_1). \end{aligned}$$

Consequently, we have

$$\phi^2 - \phi^1 \geq a \left(\frac{d^2}{D_2^2} - \frac{\varepsilon^2}{\alpha_1^2} \right) \quad \forall x \in B_\varepsilon(x_1). \quad (37)$$

It is clear that an analogous result is true by interchanging x_1 and x_2 (now with α_2 and D_1). Thus, taking $\varepsilon < \min\left(\frac{d\alpha_1}{D_2}, \frac{d\alpha_2}{D_1}\right)$ we can absorb the desired terms in the inequality and theorem 2 is proved.

3. Proof of the stability of the inverse problem

In this section, we apply the Carleman inequality of theorem 1 to the inverse problem presented in section 1. For a principal coefficient a piecewise constant and $p \in L^\infty(\Omega)$, we consider the wave equation

$$\begin{cases} u_{tt} - \operatorname{div}(a(x)\nabla u) + p(x)u = g(x, t) & \Omega \times (0, T) \\ u = h & \Gamma \times (0, T) \\ u(0) = u_0 & \Omega \\ u_t(0) = u_1 & \Omega. \end{cases} \quad (38)$$

If $g \in L^1(0, T; L^2(\Omega))$, $h \in L^2(0, T; L^2(\Gamma))$ and $u_0 \in H_0^1(\Omega)$, $u_1 \in L^2(\Omega)$, then [32, 13] equation (38) has a unique weak solution $u \in C([0, T]; H_0^1(\Omega)) \cap C^1([0, T]; L^2(\Omega))$ with continuous dependence on initial conditions and such that $\frac{\partial u}{\partial \nu} \in L^2(0, T; L^2(\Gamma))$.

In order to prove the local stability of the nonlinear application (3), that is, the problem of determining the potential p in Ω by a single measurement of the flux $a_2 \frac{\partial u}{\partial \nu}$ on Γ between $t = 0$ and $t = T$, we follow the ideas of [9] and [27]. Thus, we will first consider a linearized version of this problem, which means working on the wave equation

$$\begin{cases} y_{tt} - \operatorname{div}(a(x)\nabla y) + p(x)y = f(x)R(x, t) & \Omega \times (0, T) \\ y = 0 & \Gamma \times (0, T) \\ y(0) = 0 & \Omega \\ y_t(0) = 0 & \Omega \end{cases} \quad (39)$$

given p and R , and proving the stability of the application $f|_\Omega \mapsto \frac{\partial y}{\partial \nu}|_\Sigma$.

We will indeed prove the following result:

Theorem 3. For $x_1, x_2 \in \Omega_1$ let $R_j = \sup\{|x - y_j(x)| : x \in \Omega_2\}$, $j = 1, 2$, where y_j is defined in (7) with $x_0 = x_j$. Set $\alpha_j = d(x_j, \Gamma_1)$ and $D_0 = \max\left\{\frac{R_1 + \alpha_1}{\alpha_1}, \frac{R_2 + \alpha_2}{\alpha_2}\right\}$. With the hypothesis of theorem 1, and T, β satisfying (12), (22), (23) and (24), suppose that

- $\|p\|_{L^\infty(\Omega)} \leq m$,
- $T > D_0 \sqrt{\frac{\alpha_1}{\beta}}$,
- $R \in H^1(0, T; L^\infty(\Omega))$,
- $0 < r < |R(x, 0)|$ almost everywhere in Ω .

Then there exists $C > 0$ such that for all $f \in L^2(\Omega)$ the solution y of (39) satisfies

$$\|f\|_{L^2(\Omega)}^2 \leq C \left\| a_2 \frac{\partial y}{\partial v} \right\|_{H^1(0, T; L^2(\Gamma))}^2.$$

Proof. For each $f \in L^2(\Omega)$ and $R \in H^1(0, T; L^\infty(\Omega))$, let y be the solution of (39). We take the even extension of R and y to the interval $(-T, T)$. We call this function in the same way, and in this proof we denote by $Q = \Omega \times (-T, T)$ and $\Sigma = \Gamma \times (-T, T)$ the extended domains. Therefore, $z = y_t$ satisfies the following equation:

$$\begin{cases} z_{tt} - \operatorname{div}(a \nabla z) + pz = f(x)R_t(x, t) & Q \\ z = 0 & \Sigma \\ z(0) = 0 & \Omega \\ z_t(0) = f(x)R(x, 0) & \Omega \end{cases} \quad (40)$$

and we have the usual energy estimate

$$\|z\|_{H^1(-T, T; L^2(\Omega))} \leq C \|f R_t\|_{L^2(-T, T; L^2(\Omega))} + C \|f R(0)\|_{L^2(\Omega)}$$

that gives, since $R \in H^1(0, T; L^\infty(\Omega))$,

$$\|z\|_{H^1(-T, T; L^2(\Omega))} \leq C \|f\|_{L^2(\Omega)} \|R\|_{H^1(0, T; L^\infty(\Omega))} \leq C \|f\|_{L^2(\Omega)}. \quad (41)$$

In order to apply theorem 2 and use the appropriate Carleman estimate, we need a solution of the wave equation that vanishes at time $t = \pm T$. Thus, for $0 < \delta < T$ we take the cut-off function $\theta \in C_0^\infty(-T, T)$ such that

- $0 \leq \theta \leq 1$,
- $\theta(t) = 1$, for all $t \in (-T + \delta, T - \delta)$,

and we define $v = \theta z$. Then v satisfies

$$\begin{cases} v_{tt} - \operatorname{div}(a \nabla v) + pv = \theta(t) f(x) R_t(x, t) + 2\theta_t y_{tt} + \theta_{tt} y_t & Q \\ v = 0 & \Sigma \\ v(0) = 0 & \Omega \\ v_t(0) = f(x) R(x, 0) & \Omega \\ v(\pm T) = v_t(\pm T) = 0 & \Omega. \end{cases} \quad (42)$$

Take $j \in \{1, 2\}$, and let y be the function defined in (7) and ϕ the weight function, corresponding to the point $x_j \in \Omega_1$. Note that

$$\phi(x, t) \leq \phi(x, 0) \quad \forall (x, t) \in (0, T) \times \Omega. \quad (43)$$

Moreover, by the definition of ρ and \bar{a} (see (8) and the definitions below) we also have

$$\frac{|x - x_j|}{\rho(x)} \leq 1 + \frac{|x - y(x)|}{\rho(x)} \leq 1 + \frac{R_j}{\alpha_j} \leq D_0$$

and then

$$\begin{aligned}\phi(x, t) &= \bar{a} \frac{|x - x_j|^2}{\rho^2(x)} - \beta t^2 + M \\ &\leq \bar{a} D_0^2 - \beta t^2 + M.\end{aligned}\quad (44)$$

Then, by the choice of $T > D_0 \sqrt{\frac{a_1}{\beta}}$ we get

$$\phi(x, \pm T) < M \leq \phi(x, 0) \quad \forall x \in \Omega. \quad (45)$$

Thus, taking δ small enough, it is also true that

$$\phi(x, t) < M \leq \phi(x, 0), \quad (46)$$

for all $x \in \Omega$ and $t \in [-T, -T + \delta] \cup [T - \delta, T]$.

From now on, $C > 0$ will denote a generic constant depending on $\Omega, T, \beta, \theta, x_1, x_2, \delta, s_0$ and λ_0 but independent of $s > s_0$ and $\lambda > \lambda_0$. We will occasionally use the notation ∂_t for the time derivative.

As in the proof of theorem 2, we set $\varphi = e^{\lambda\phi}$, $w_j = e^{s\varphi} v_j$ and

$$P_1(w) = w_{tt} - a\Delta w + s^2\lambda^2\varphi^2 E(\phi)w.$$

It is easy to check that

$$\sum_{j=1}^2 \langle P_1 w_j, \partial_t w_j \rangle_{L^2(\Omega_j \times (0, T))} = \frac{1}{2} \int_{\Omega} |\partial_t w(0)|^2 + X \quad (47)$$

where X is the sum of negligible terms such that

$$\sum_{j=1}^2 \langle P_1 w_j, \partial_t w_j \rangle_{L^2(\Omega_j \times (0, T))} \geq \frac{1}{2} \int_{\Omega} |\partial_t w(0)|^2 - Cs^2\lambda^3 \int_0^T \int_{\Omega} \varphi^3 |w|^2.$$

Since we have $w_t(0) = e^{s\varphi(0)} v_t(0) = e^{s\varphi(0)} f(x) R(x, 0)$ and $|R(x, 0)| \geq r$, we get (recall that $Q = \Omega \times (-T, T)$, $Q_j = \Omega_j \times (-T, T)$ and so on)

$$r^2 \int_{\Omega} e^{2s\varphi(0)} |f|^2 \leq C \left(\sum_{j=1}^2 \langle P_1 w_j, \partial_t w_j \rangle_{L^2(\Omega_j \times (0, T))} + s^2\lambda^3 \iint_Q \varphi^3 |w|^2 \right).$$

In order to apply the Carleman estimate (corollary 22) we consider both weight functions φ^1 and φ^2 , corresponding to x_1 and $x_2 \in \Omega_1$, and we apply the previous estimates to $w_j^k = e^{s\varphi^k} v_j$ for $j, k = 1, 2$ and sum up the inequalities. We obtain, for $s > s_0$ and $\lambda > \lambda_0$, using Cauchy–Schwarz inequality, the following:

$$\begin{aligned}r^2 \int_{\Omega} (e^{2s\varphi^1(0)} + e^{2s\varphi^2(0)}) |f|^2 &\leq C \sum_{k=1}^2 \left(\sum_{j=1}^2 \langle P_1^{\phi^k} w_j^k, \partial_t w_j^k \rangle_{L^2(\Omega_j \times (0, T))} + s^2\lambda^3 \iint_Q (\varphi^k)^3 |w^k|^2 \right) \\ &\leq C \sum_{j,k=1}^2 \left(\frac{1}{\sqrt{s}} |P_1^{\phi^k} w_j^k|_{L^2(Q_j)}^2 + \sqrt{s} |\partial_t w_j^k|_{L^2(Q_j)}^2 \right) + Cs^2\lambda^3 \sum_{k=1}^2 \iint_Q (\varphi^k)^3 |w^k|^2 \\ &\leq \frac{C}{\sqrt{s}} \sum_{j,k=1}^2 (|P_1^{\phi^k} w_j^k|_{L^2(Q_j)}^2 + \|w_j^k\|_{Q_j, \varphi^k}^2).\end{aligned}$$

Now, applying corollary 22, we get

$$\begin{aligned} r^2 \int_{\Omega} (e^{2s\varphi^1(0)} + e^{2s\varphi^2(0)}) |f|^2 &\leq \frac{C}{\sqrt{s}} \sum_{k=1}^2 (|P_1^{\phi^k} w^k|_{L^2(Q_0)}^2 + \|w^k\|_{Q_0, \varphi^k}^2) \\ &\leq \frac{C}{\sqrt{s}} \sum_{k=1}^2 \left(|e^{s\varphi^k} L_P v|_{L^2(Q_0)}^2 + s\lambda \int \int_{\Sigma_+^{\phi^k}} \varphi^k \left| a_2 \frac{\partial w^k}{\partial \nu} \right|^2 \right). \end{aligned}$$

On the one hand, since we have $\theta_t = 0$ in $[-T + \delta, T - \delta]$, then from estimates (41), (43) and (46), we obtain, $\forall k = 1, 2$,

$$\begin{aligned} \int \int_{Q_0} e^{2s\varphi^k} |L_P v|^2 &= \int \int_{Q_0} e^{2s\varphi^k} |\theta f R_t + 2\theta_t y_{tt} + \theta_{tt} y_t|^2 \\ &\leq C \int \int_{Q_0} e^{2s\varphi^k} |f|^2 |R_t|^2 + C \int \int_{Q_0} e^{2s\varphi^k} (|\theta_t z_t|^2 + |\theta_{tt} z|^2) \\ &\leq C \int \int_{Q_0} e^{2s\varphi^k(0)} |f|^2 |R_t|^2 + C \left(\int_{-T}^{-T+\delta} + \int_{T-\delta}^T \right) \int_{\Omega_0} e^{2se^{\lambda M}} (|z_t|^2 + |z|^2) \\ &\leq C \|R\|_{H^1(0, T; L^\infty)} \int_{\Omega} e^{2s\varphi^k(0)} |f|^2 + C e^{2se^{\lambda M}} \|z\|_{H^1(-T, T; L^2(\Omega))}^2 \\ &\leq C \int_{\Omega} e^{2s\varphi^k(0)} |f|^2 + C e^{2se^{\lambda M}} \|f\|_{L^2(\Omega)} \\ &\leq C \int_{\Omega} e^{2s\varphi^k(0)} |f|^2. \end{aligned}$$

Now, recalling the notation stated in (15), we have $\Sigma_+^{\phi^1} \cup \Sigma_+^{\phi^2} \subset \Sigma$ and $|\frac{\partial w^k}{\partial \nu}| = e^{s\varphi^k} |\frac{\partial v}{\partial \nu}|$ on Σ for each k , so we finally obtain

$$\begin{aligned} r^2 \int_{\Omega} (e^{2s\varphi^1(0)} + e^{2s\varphi^2(0)}) |f|^2 &\leq \frac{C}{\sqrt{s}} \int_{\Omega} (e^{2s\varphi^1(0)} + e^{2s\varphi^2(0)}) |f|^2 \\ &\quad + C \sqrt{s} \lambda \int \int_{\Sigma} (\varphi^1 e^{2s\varphi^1} + \varphi^2 e^{2s\varphi^2}) \left| a_2 \frac{\partial v}{\partial \nu} \right|^2. \end{aligned} \quad (48)$$

For s large enough, the left-hand side in (48) can absorb the first term of the right-hand side. Therefore, since φ^k and θ are bounded on Σ and $\frac{\partial z}{\partial \nu}$ is an even function with respect to $t \in [-T, T]$, we obtain

$$\begin{aligned} \int_{\Omega} |f|^2 &\leq C \int \int_{\Sigma} (\varphi^1 e^{2s\varphi^1} + \varphi^2 e^{2s\varphi^2}) \left| a_2 \frac{\partial v}{\partial \nu} \right|^2 \\ &\leq C \int \int_{\Sigma} \left| a_2 \frac{\partial z}{\partial \nu} \right|^2 \\ &= 2C \int_0^T \int_{\Gamma} \left| a_2 \frac{\partial z}{\partial \nu} \right|^2. \end{aligned}$$

and this ends the proof of theorem 3. \square

We will end this paper by the proof of theorem 1 which is a direct consequence of theorem 3. Indeed, if we set $\tilde{y} = u(q) - u(p)$, $f = p - q$ and $R = u(p)$, then \tilde{y} is the solution of

$$\begin{cases} \tilde{y}_{tt} - \operatorname{div}(a\nabla\tilde{y}) + (p - f)\tilde{y} = f(x)R(x, t) & (0, T) \times \Omega \\ \tilde{y} = 0 & (0, T) \times \Sigma \\ \tilde{y}(0) = 0 & \Omega \\ \tilde{y}_t(0) = 0 & \Omega \end{cases} \quad (49)$$

where $q = p - f \in \mathcal{U}$, with \mathcal{U} bounded in $L^\infty(\Omega)$ from the hypothesis of theorem 1. The key point is that in the proof of theorem 3 all the constants $C > 0$ depend on the L^∞ -norm of the potential as stated in corollary 2. Thus, with $q \in \mathcal{U}$, we are actually, with equation (49), in a situation similar to the linear inverse problem related to equation (39) and we then obtain the desired result.

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See endnote 1

See endnote 2

Endnotes

- (1) Author: Please update reference [17].
- (2) Author: Please provide the surname of the author in references [37, 38].

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